## GR Assignments 01



## 1. Free Relativistic Particle in Arbitrary Coordinates

As you saw in the lecture, transforming from inertial (Minkowski) coordinates $\xi^{a}$ to arbitrary coordinates $x^{\mu}$, the free-particle equation of motion $d^{2} \xi^{a} / d \tau^{2}=0$ becomes

$$
\begin{equation*}
\frac{d^{2} x^{\mu}}{d \tau^{2}}+\Gamma_{\nu \lambda}^{\mu} \frac{d x^{\nu}}{d \tau} \frac{d x^{\lambda}}{d \tau}=0 \tag{1}
\end{equation*}
$$

where

$$
\begin{equation*}
\Gamma_{\nu \lambda}^{\mu}=\frac{\partial x^{\mu}}{\partial \xi^{a}} \frac{\partial^{2} \xi^{a}}{\partial x^{\nu} \partial x^{\lambda}} \equiv J_{a}^{\mu} J_{\nu \lambda}^{a} \tag{2}
\end{equation*}
$$

Show that this pseudo-force term $\Gamma$ arising in the $x^{\mu}$ coordinate system is related to the metric

$$
\begin{equation*}
g_{\mu \nu}=J_{\mu}^{a} J_{\nu}^{b} \eta_{a b} \tag{3}
\end{equation*}
$$

in these coordinates by

$$
\begin{align*}
\Gamma_{\nu \lambda}^{\mu} & =g^{\mu \rho} \Gamma_{\rho \nu \lambda} \\
\Gamma_{\mu \nu \lambda} & =\frac{1}{2}\left(g_{\mu \nu}, \lambda+g_{\mu \lambda, \nu}-g_{\nu \lambda, \mu}\right) \tag{4}
\end{align*}
$$

where $g^{\mu \nu}$ is the inverse metric and $g_{\mu \nu},_{\lambda}$ is short-hand for the partial derivative $\left(\partial g_{\mu \nu} / \partial x^{\lambda}\right)$.

## 2. GEODESICS

(a) Geodesics and Euler-Lagrange Equations: Show that the Euler-Lagrange equations

$$
\begin{equation*}
\frac{d}{d \tau} \frac{\partial \mathcal{L}}{\partial \dot{x}^{\mu}}-\frac{\partial \mathcal{L}}{\partial x^{\mu}}=0 \tag{5}
\end{equation*}
$$

for the Lagrangian

$$
\begin{equation*}
\mathcal{L}=\frac{1}{2} g_{\mu \nu} \frac{d x^{\mu}}{d \tau} \frac{d x^{\nu}}{d \tau} \tag{6}
\end{equation*}
$$

are the geodesic equations

$$
\begin{equation*}
\frac{d^{2} x^{\mu}}{d \tau^{2}}+\Gamma_{\nu \lambda}^{\mu} \frac{d x^{\nu}}{d \tau} \frac{d x^{\lambda}}{d \tau}=0 \tag{7}
\end{equation*}
$$

where the Christoffel symbols $\Gamma_{\nu \lambda}^{\mu}$ associated to a metric $g_{\mu \nu}$ are defined as in Exercise 1 (eq. 4), but now for an arbitrary metric.
(b) $\mathcal{L}$ is a constant of motion: Show that $\mathcal{L}$ is constant along any geodesic, i.e. that

$$
\begin{equation*}
\frac{d}{d \tau}\left(g_{\mu \nu} \frac{d x^{\mu}}{d \tau} \frac{d x^{\nu}}{d \tau}\right)=0 \tag{8}
\end{equation*}
$$

for $x^{\mu}(\tau)$ a solution of the geodesic equation.
(c) Geodesics on the Two-Sphere $S^{2}$ : The metric of a 2 -sphere with radius $R$ is

$$
\begin{equation*}
d s^{2}=R^{2}\left(d \theta^{2}+\sin ^{2} \theta d \phi^{2}\right) . \tag{9}
\end{equation*}
$$

Calculate all its Christoffel symbols, show that the geodesic equations agree with the Euler-Lagrange equations of the Lagrangian

$$
\begin{equation*}
\mathcal{L}=\frac{1}{2}\left(\dot{\theta}^{2}+\sin ^{2} \theta \dot{\phi}^{2}\right), \tag{10}
\end{equation*}
$$

and show that the great circles (longitudes) $(\theta(\tau), \phi(\tau))=\left(\tau, \phi_{0}\right)$ satisfy the geodesic equation.

