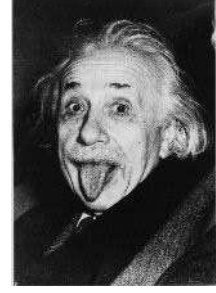


FORMELSAMMLUNG KLASSISCHE FELDTHEORIE



8 KLASSISCHE FELDTHEORIE (IM ENGEREN SINN)

	Klassische Mechanik	Klassische Feldtheorie
unabhängige Variable	Zeit t oder Eigenzeit τ	Raum(-Zeit) Koordinaten x^α $\alpha = 0, \dots, d-1$ oder $\alpha = 1, \dots, d$
abhängige Variable	Koordinaten $q^i(t)$ oder $x^\alpha(\tau)$	Felder $\Phi^A(x^\alpha)$ Φ^A : Skalar-, Vektor-, Tensorfelder etc.
Ableitungen	$(d/dt)q^i(t) = \dot{q}^i(t)$ oder $\dot{x}^\alpha(\tau)$	partielle Ableitungen $\partial_\alpha \Phi^A(x)$
Lagrange-Funktion	$L = L(q^i, \dot{q}^i; t)$	$L = L(\Phi^A, \partial_\alpha \Phi^A; x^\alpha)$
Wirkung	$S[q] = \int dt L$	$S[\Phi] = \int d^d x L$
Variation	$q^i(t) \rightarrow q^i(t) + \delta q^i(t)$	$\Phi^A(x) \rightarrow \Phi^A(x) + \delta \Phi^A(x)$
Euler-Lagrange Gleichungen	$(\partial L / \partial q^i) - (d/dt)(\partial L / \partial \dot{q}^i) = 0$ $\frac{d}{dt} = \frac{\partial}{\partial t} + \dot{q}^i(t) \frac{\partial}{\partial q^i(t)} + \dots$	$(\partial L / \partial \Phi^A) - (d/dx^\alpha)(\partial L / \partial (\partial_\alpha \Phi^A)) = 0$ $\frac{d}{dx^\alpha} := \frac{\partial}{\partial x^\alpha} + \partial_\alpha \Phi^A(x) \frac{\partial}{\partial \Phi^A(x)} + \dots$

1. Reelles Skalarfeld Φ im Minkowski-Raum

(a) Freies masseloses Feld

$$\begin{aligned}
 L(\Phi, \partial_\alpha \Phi) &= -\frac{1}{2} \eta^{\alpha\beta} \partial_\alpha \Phi \partial_\beta \Phi \\
 S[\Phi] &= \int d^4 x L(\Phi, \partial_\alpha \Phi) = -\frac{1}{2} \int d^4 x \eta^{\alpha\beta} \partial_\alpha \Phi \partial_\beta \Phi \\
 \delta S[\Phi] &= 0 \quad \Rightarrow \quad \square \Phi \equiv \eta^{\alpha\beta} \partial_\alpha \partial_\beta \Phi = 0
 \end{aligned} \tag{1}$$

(b) Freies massives Feld

$$\begin{aligned}
 L(\Phi, \partial_\alpha \Phi) &= -\frac{1}{2} \eta^{\alpha\beta} \partial_\alpha \Phi \partial_\beta \Phi - \frac{1}{2} m^2 \Phi^2 \\
 S[\Phi] &= \int d^4 x L(\Phi, \partial_\alpha \Phi) = -\frac{1}{2} \int d^4 x \left(\eta^{\alpha\beta} \partial_\alpha \Phi \partial_\beta \Phi + m^2 \Phi^2 \right) \\
 \delta S[\Phi] &= 0 \quad \Rightarrow \quad (\square - m^2) \Phi = 0
 \end{aligned} \tag{2}$$

(c) Feld mit (Selbst-)Wechselwirkung

$$\begin{aligned}
 L(\Phi, \partial_\alpha \Phi) &= -\frac{1}{2} \eta^{\alpha\beta} \partial_\alpha \Phi \partial_\beta \Phi - V(\Phi) \\
 S[\Phi] &= \int d^4 x L(\Phi, \partial_\alpha \Phi) = \int d^4 x \left(-\frac{1}{2} \eta^{\alpha\beta} \partial_\alpha \Phi \partial_\beta \Phi - V(\Phi) \right) \\
 \delta S[\Phi] &= 0 \quad \Rightarrow \quad \square \Phi = V'(\Phi) \equiv \frac{\partial V(\Phi)}{\partial \Phi}
 \end{aligned} \tag{3}$$

2. Maxwell-Theorie

(a) Vakuum-Maxwellgleichungen

$$\begin{aligned}L_0(A_\beta, \partial_\alpha A_\beta) &= -\frac{1}{4} F_{\alpha\beta} F^{\alpha\beta} \\S_0[A] &= \int d^4x L_0(A_\beta, \partial_\alpha A_\beta) = -\frac{1}{4} \int d^4x F_{\alpha\beta} F^{\alpha\beta} \\ \delta S_0[A] = 0 &\Rightarrow \partial_\alpha F^{\alpha\beta} = 0\end{aligned}\tag{4}$$

(b) Kopplung an externen Strom J^α (mit $\mu_0 = 1$)

$$\begin{aligned}L(A_\beta, \partial_\alpha A_\beta, J) &= L_0(A_\beta, \partial_\alpha A_\beta) + A_\alpha J^\alpha \\S[A, J] &= \int d^4x L(A_\beta, \partial_\alpha A_\beta, J) = S_0[A] + \int d^4x A_\alpha J^\alpha \equiv S_0[A] + S_I[A, J] \\ \delta S[A, J] = 0 &\Rightarrow \partial_\alpha F^{\alpha\beta} = -J^\beta\end{aligned}\tag{5}$$

(c) Eichinvarianz der Wirkung

$$\begin{aligned}S_0[A_\alpha + \partial_\alpha \Psi] &= S_0[A_\alpha] \\S_I[A_\alpha + \partial_\alpha \Psi, J] &= S_I[A_\alpha, J] \quad (\text{mod Randterme}) \Leftrightarrow \partial_\alpha J^\alpha = 0\end{aligned}\tag{6}$$